

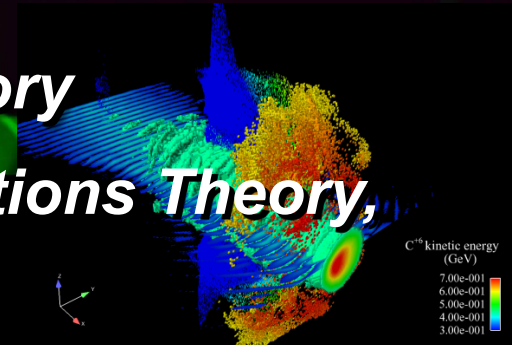
**The little foil that can:
ion-, electron- and photon-beams from laser-
nanofoil interactions**

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Los Alamos National Laboratory

*Frontiers in Intense Laser-Matter Interactions Theory,
Garching, 3. March. 2010*

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Excellent.**

Outline

Nanofolios can be ...

- ... be a GeV ion source
- ... do laser pulse shaping
- ... be a coherent x-ray source
- Summary

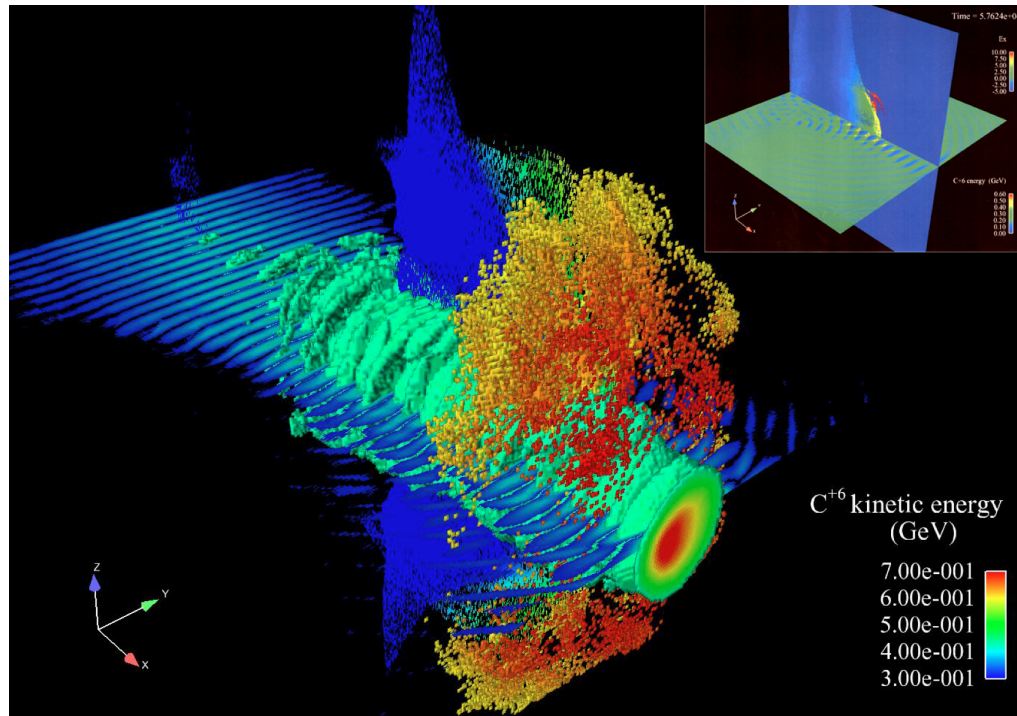
Nanotargets can be an ion source

- Break-Out Afterburner (BOA) and Radiation Pressure Acceleration (RPA)
 - 0.5 GeV peak energies
 - High conversion efficiency
 - Quasi-monoenergetic ions

Laser – Nanotarget Interaction:

Targets 30nm, 58nm Diamond-Like Carbon foils @ 821 n_{cr}

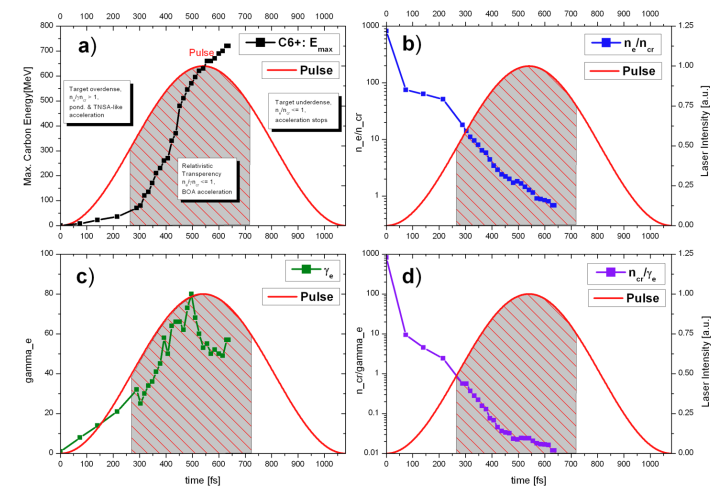
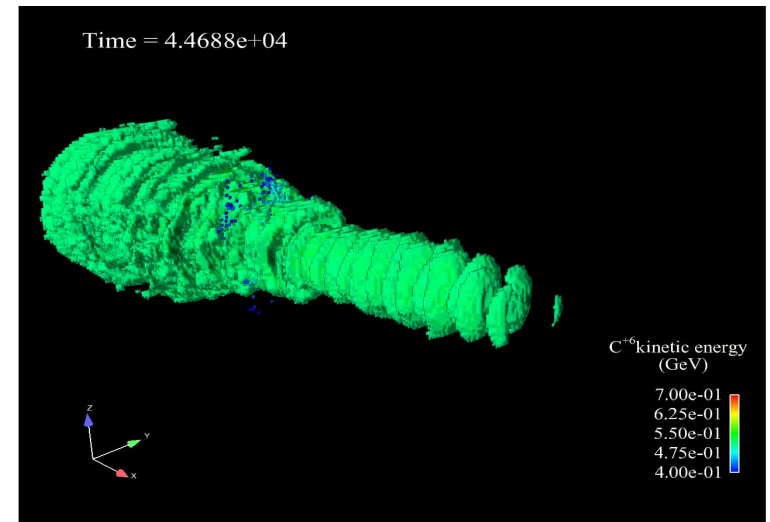
Lin. Pol.



3D-VPIC simulation of a 10^{21} W/cm², 304 fs, laser pulse interacting with a 30nm DLC foil

Parameters of BOA phase:

- Duration: ~500fs
- Spatial extend: ~15 μ m



2D-VPIC simulation: 10^{21} W/cm², 540 fs, laser interacting with a 58nm DLC foil

Energy transfer: laser \rightarrow electrons \rightarrow ions

Dispersion relation consistent with rel. Buneman instability

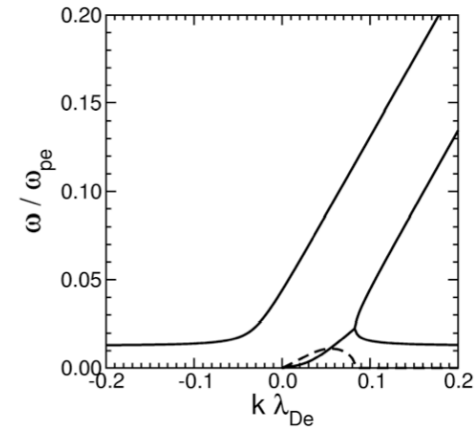
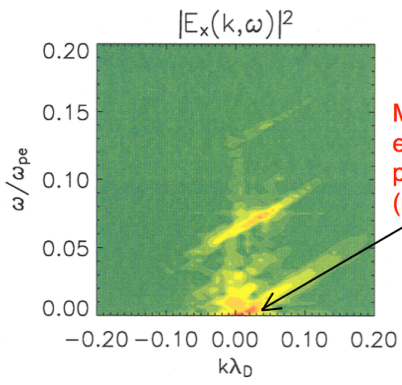
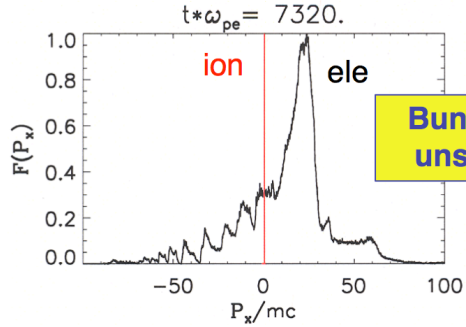
Parameters of BOA phase:

- Duration: ~ 500 fs
- Spatial extend: $\sim 15 \mu\text{m}$

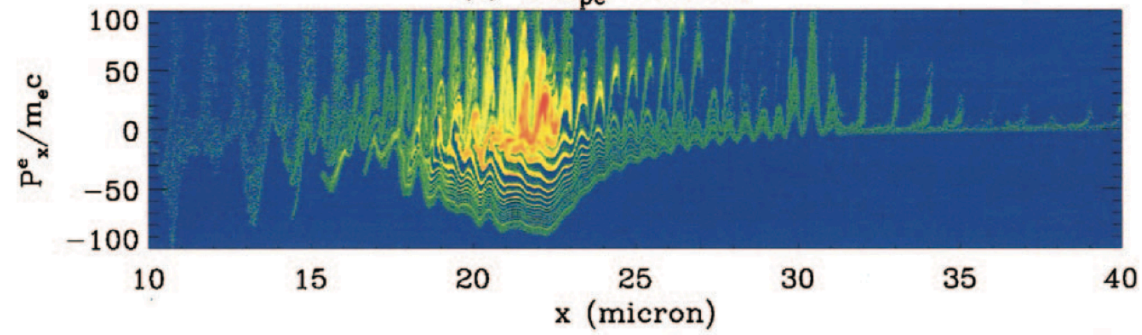
Parameters of instability:

- Growth rate: 1.7 fs
- Phase speed: 0.1 c
- Spatial extend: ~ 50 -500 nm

Electron distribution function at peak of E_x field & bulk of C ions



(a) $t * \omega_{pe} = 7320$.



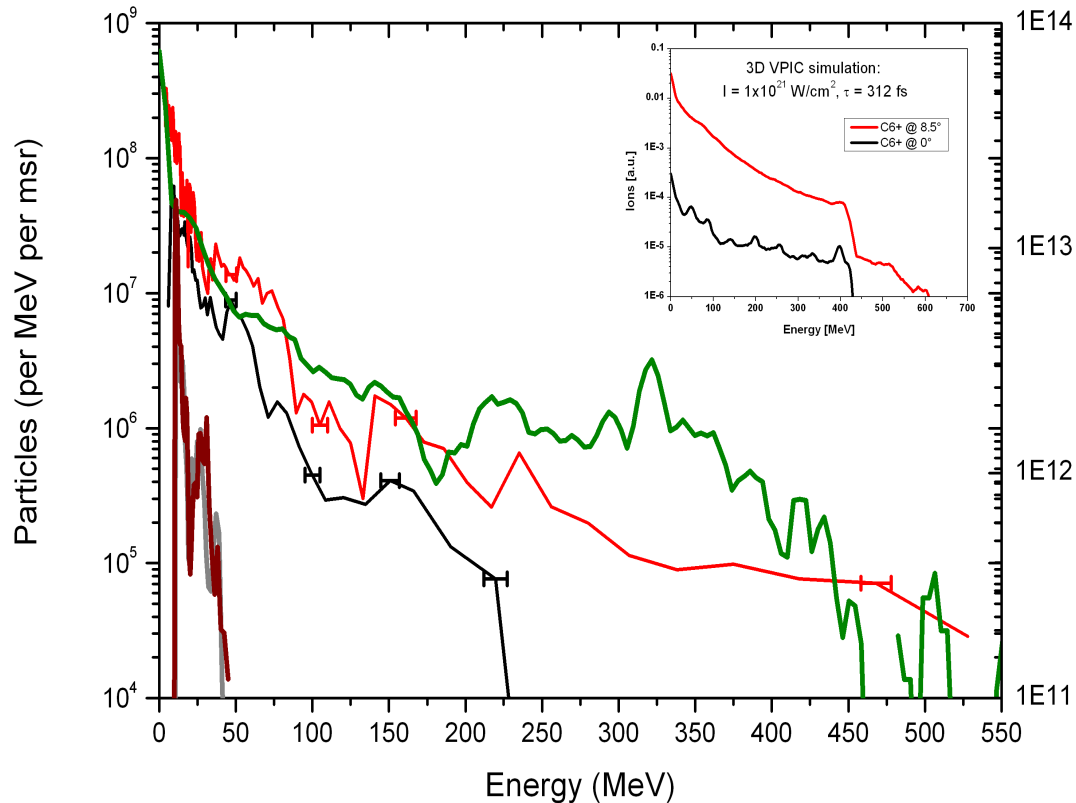
Each half-cycle of the laser launches high- γ electron bunch
 \Rightarrow many independent occurrences of instability

BOA: Comparing Experiment & Sim.

58nm DLC foil, $I_{\text{peak}} = 5 \times 10^{20} \text{ W/cm}^2$, $I_{\text{ave}} = 2 \times 10^{20} \text{ W/cm}^2$, $t = 540 \text{ fs}$,



Lin. Pol.



- Simulations agrees with measured spectra (energy, number, angular distribution)
- Protons have same velocity as Carbon
- Spectra retain mono-energetic remnant from adiabatic phase

Theory: reduced model for max. ion energy

Yan, Tajima, Hegelich, Appl. Phys. B. (2009) in press, (arXiv:0904.1466)

Lin. Pol. 

$$E_{\text{sum}} = E_{\text{max},i} + E_{\text{max},BOA}$$

$$\varepsilon_{\text{max},i}(t_1) = (2\alpha + 1)Q\bar{E}_0(t_1)((1 + \omega t_1)^{1/2\alpha+1} - 1)$$

$$\bar{E}_0(t_1) = m_e c^2 \int_0^{t_1} (\sqrt{T(t')a^2(t') + 1} - 1) \frac{dt'}{t_1}$$

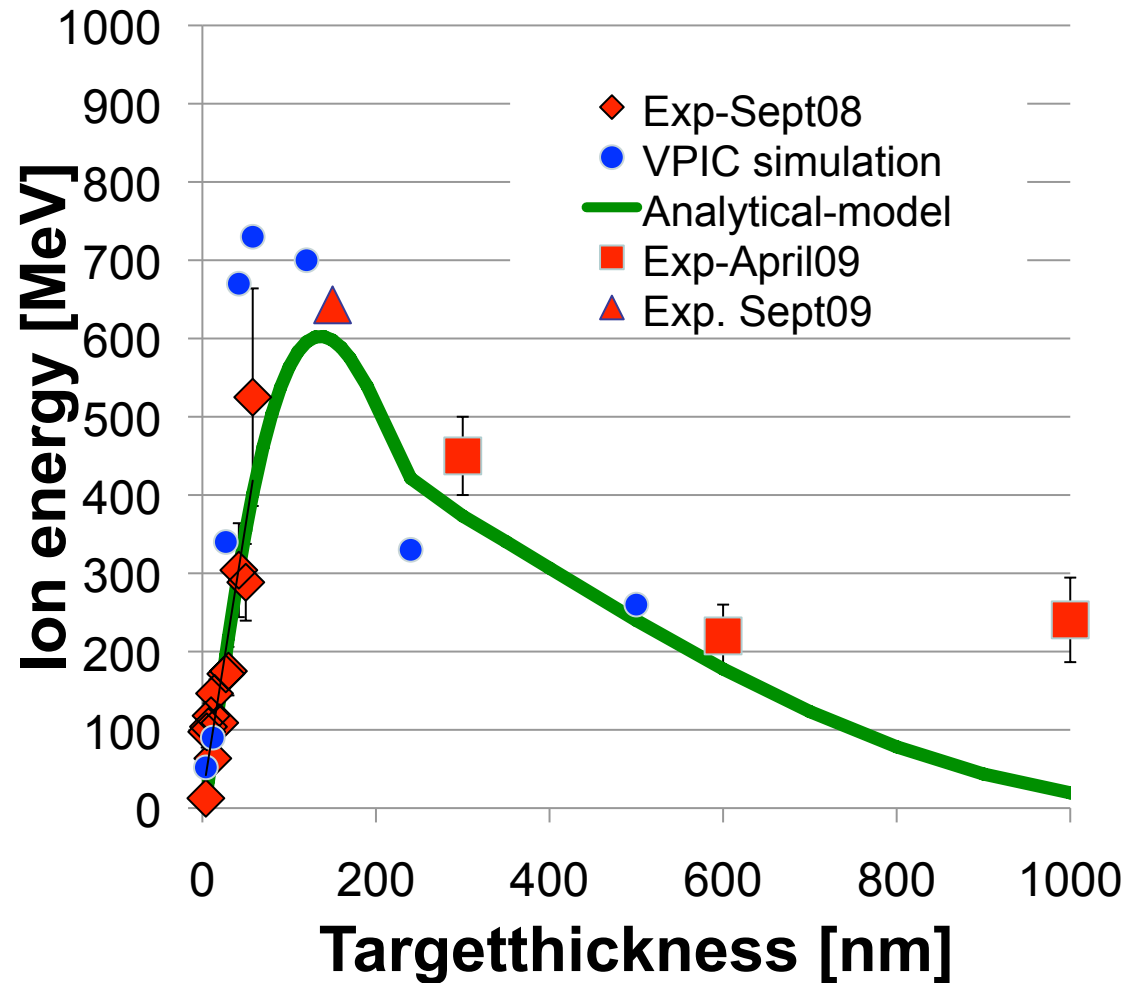
$$\varepsilon_{\text{max},i,BOA} = (2\alpha + 1)Q\bar{E}_0((1 + \omega(t_2 - t_1))^{1/2\alpha+1} - 1)$$

$$t_1 : n_e / \gamma n_{cr} \sim 1$$

$$t_1 = \left(\frac{M}{m} \frac{3N^2 d^2}{q_i \omega^2 c^2 a_0^3} \right)^{1/4} = \left(\frac{12}{q_i \pi^2} \frac{M}{m} \frac{N^2 d^2 \tau^2}{a_0^3 c^2} \right)^{1/4}$$

$$t_2 = t_1 + \Delta t : n_e / n_{cr} \sim 1$$

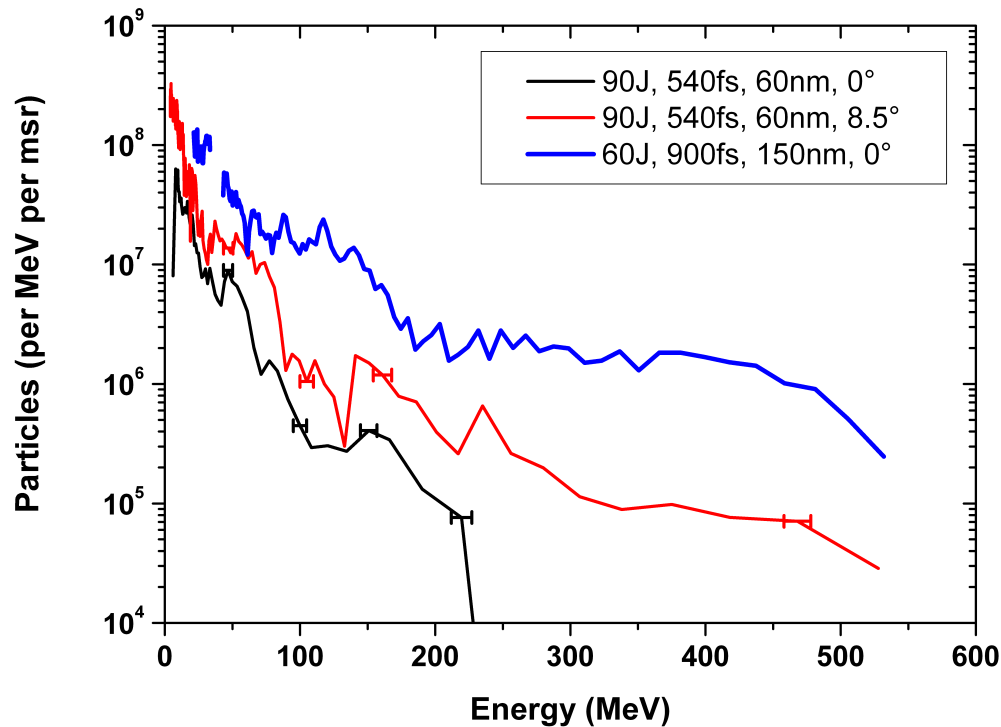
$$\Delta t = \frac{x_2 - x_1}{C_{s,\text{max}}} \frac{1}{\sin(\omega t_1)} = \frac{x_1(\gamma^{1/3} - 1)}{C_s} \frac{1}{\sin(\omega t_1)} = \frac{Nd(\gamma^{1/3} - 1)}{\gamma C_s} \frac{1}{\sin(\omega t_1)}$$



BOA: Ongoing Opt. of laser-target param.

150nm DLC foil, $I_{\text{peak}} = 2.5 \times 10^{20} \text{ W/cm}^2$, $I_{\text{ave}} = 8 \times 10^{19} \text{ W/cm}^2$

Lin. Pol. 

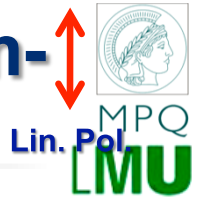


- Further contrast improvement on 5-10ps scale: closer to ideal pulse shape
- Target thickness optimized to laser parameters
- 2x improvement of energy @ same angle
- 2 orders of magnitude higher efficiency



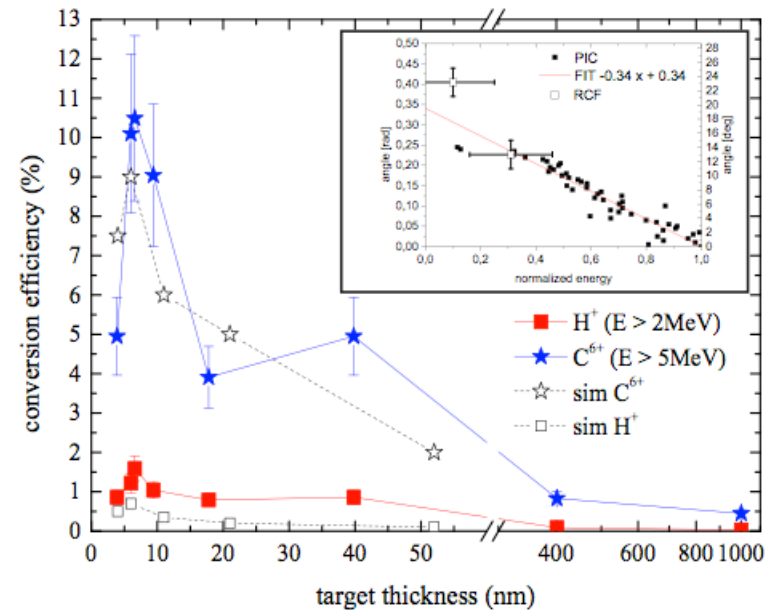
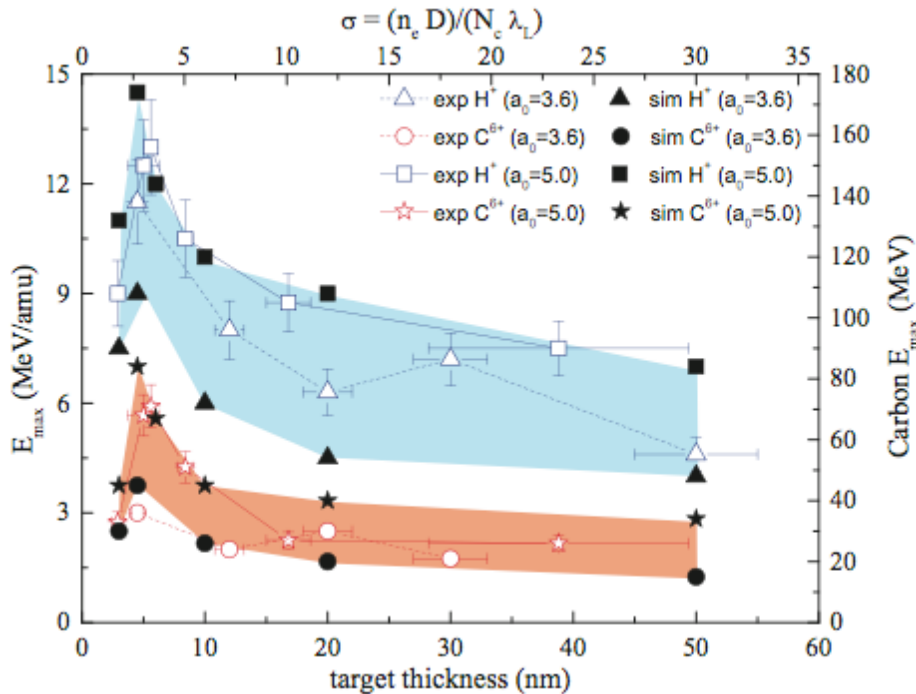
Max-Born-Institut

Enhanced Efficiency and Energy from nm-targets: Steinke et al., AP B (2010)



Emax (Proton): 14 MeV
Emax (Carbon): 80 MeV

CE (Proton): 2%
CE (Carbon): 11%

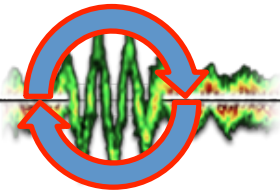


- 10Hz Ti:Sapphire laser system: 700 mJ, 45 fs, $3 \times 10^{19} \text{ W/cm}^2$
- double plasma mirror for contrast enhancement
- DLC-foils of thickness 50nm down to sub-5nm (shot at normal incidence)

Experiments @ MBI Berlin:
 A. Henig, Habs, M. Hegelich, R. Hörlein, D. Kiefer, D. Jung, X. Yan, D. (LMU)
 S. Steinke, T. Sokollik, M. Schnürer, P. Nickles, W. Sandner (MBI)



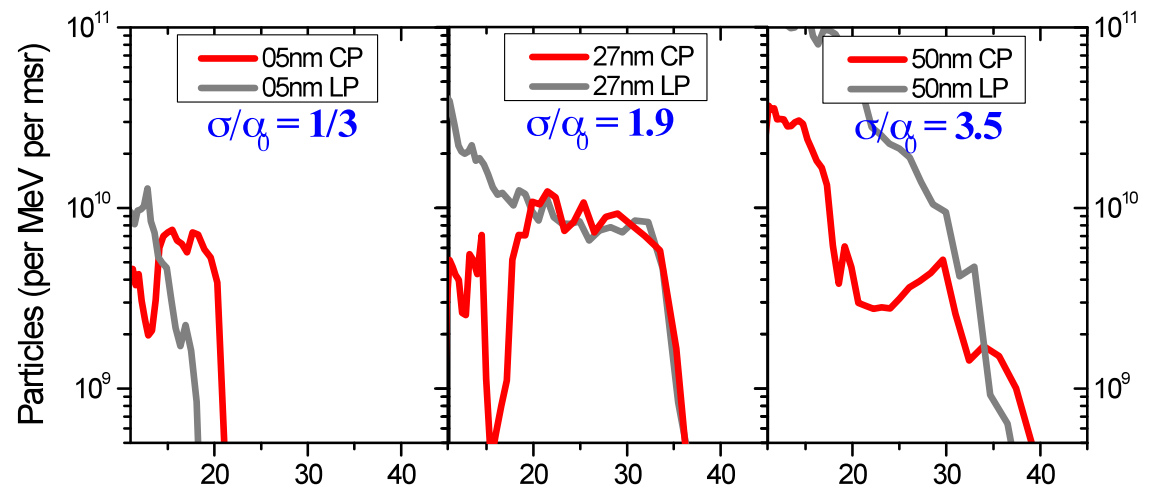
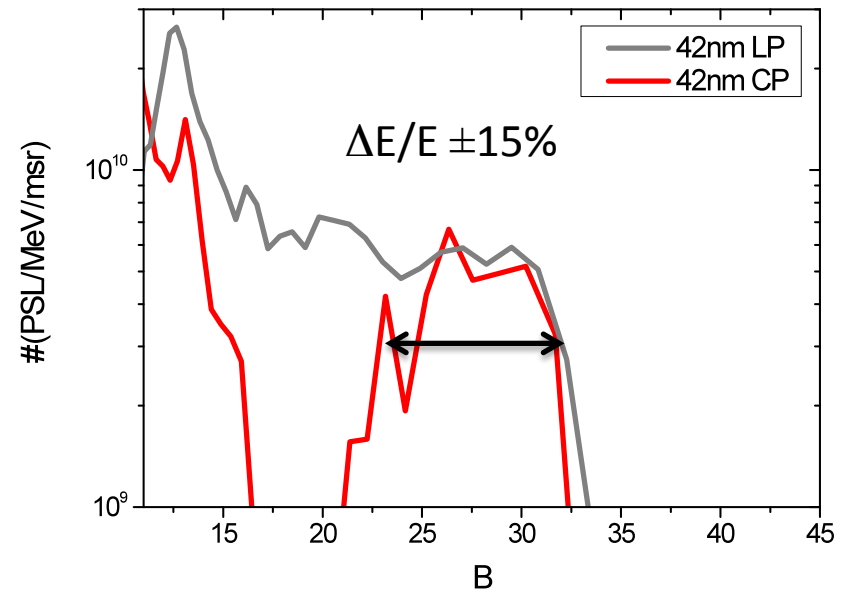
Proton data from experiment using CP



Protons

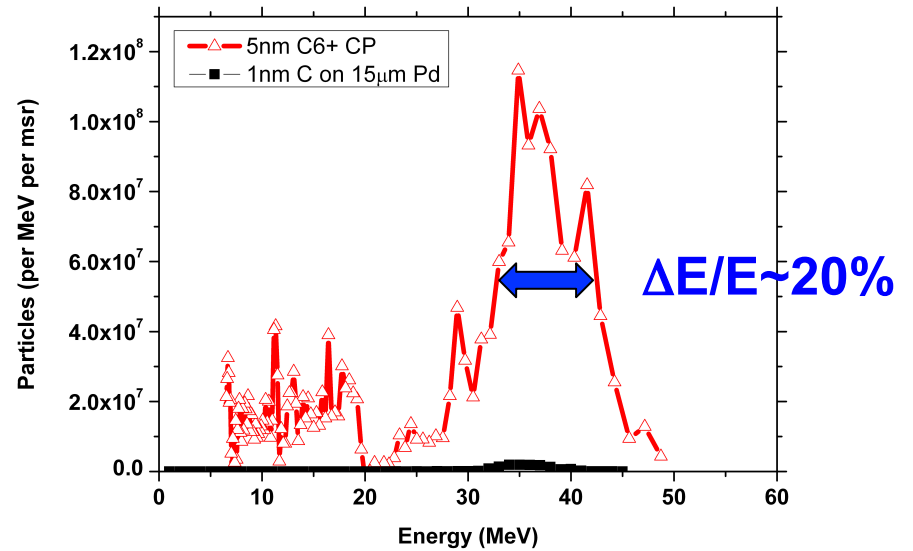
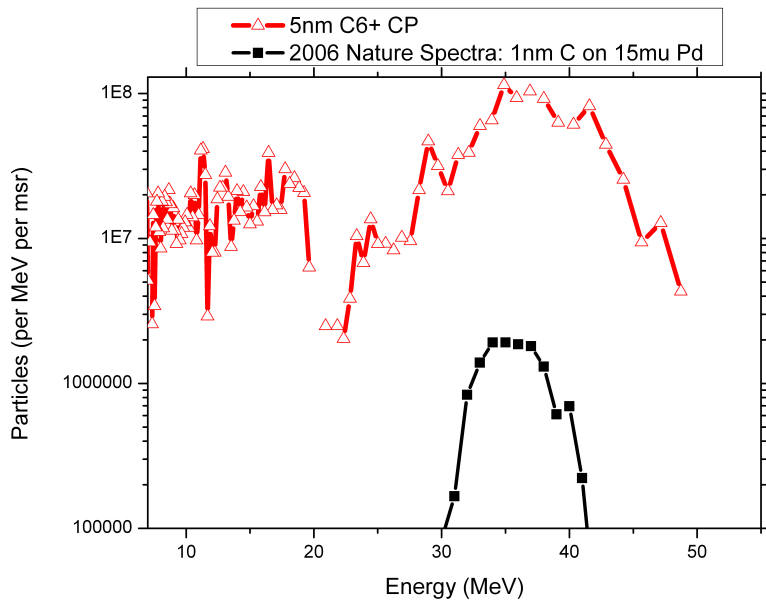
σ Scan from 3 to 50 ($a_0=12$)

- Equals thickness of 3nm to 60nm
- For $1/3 \lesssim \sigma/a_0 \lesssim 3$ we see monoenergetic proton beams with CP, otherwise exponential spectra
- First experimental generation of monoenergetic proton beams with circular polarized light showing signatures of light pressure acceleration (RPA/PSA)
- Monoenergetic spectra measured 1.2m behind the target; beam quasi-neutrality indicated by theory



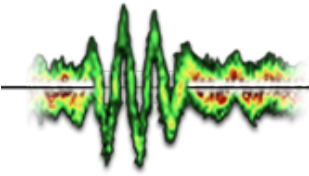
Trident: Carbon Acceleration w Circ. Pol.

typical pulse parameter: 80J, 550fs, 3×10^{20} W/cm², contrast < 10^{-11} @ 50ps



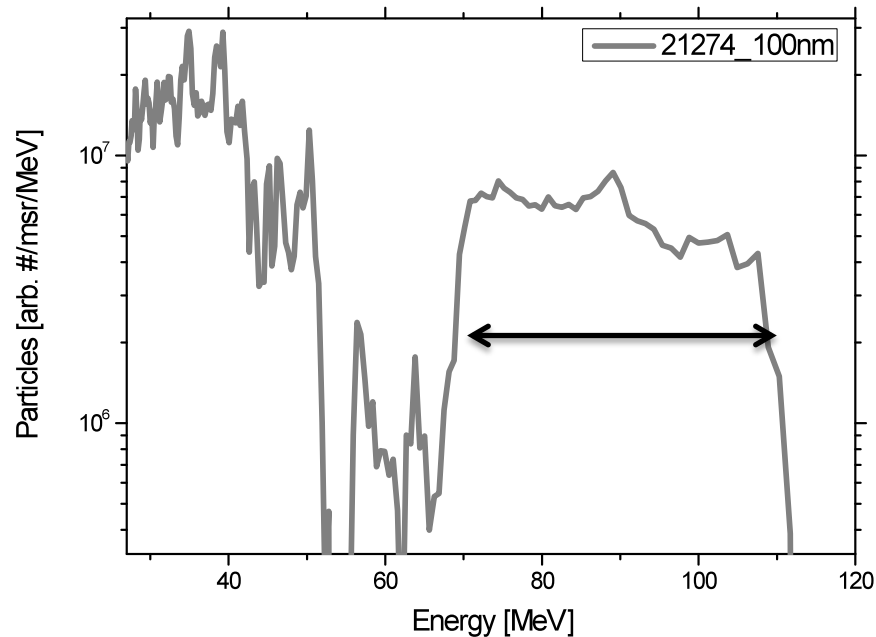
Almost 2 orders of magnitude more particles into peak than in the 2006 Nature experiment.

Approximately 20% energy spread with far from optimal laser parameters.



Main results from F/8 OAP and LP

$\Delta E/E \pm 20\% @ 90\text{MeV}$



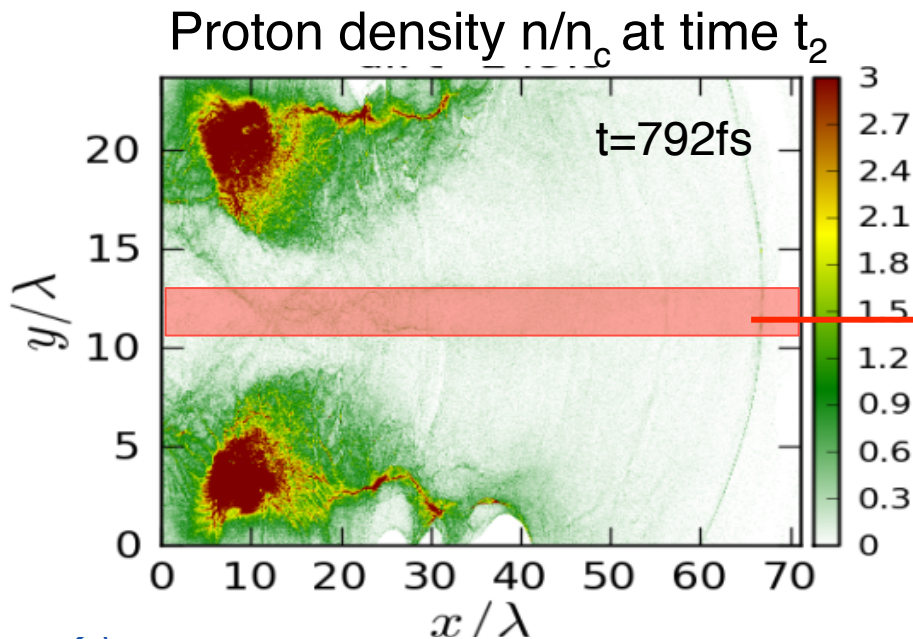
- 1/10th of intensity due to increased spot size $\sim 20 \mu\text{m}$ decreased $a_0 \sim 3$
- linear instead of circular polarization
- mono-energetic feature at 90 MeV
- increased energy spread of $\Delta E/E \sim 40\%$
- Mono-energetic feature is reproducible with all thicknesses from 5 nm to 100 nm

H target

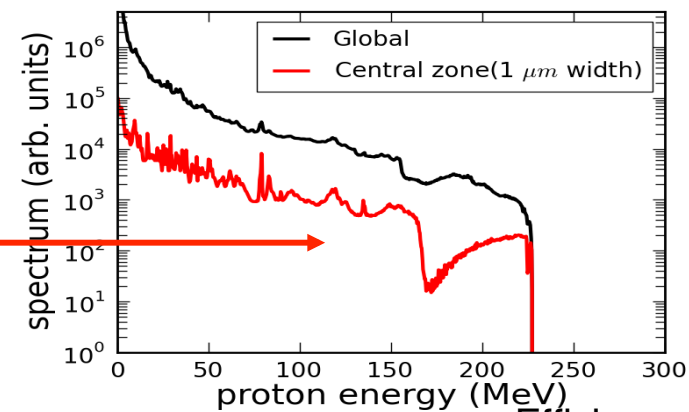
Cyrogenic hydrogen target leads to higher energy and efficiency than CH₂

Cyrogenic hydrogen target at (0.07g/cm³), $n_e=42.6 n_{cr}$

Target Thickness l	400 nm	1000 nm	1500nm	2500 nm
H ⁺ E _{max}	145 MeV	205 MeV	230MeV	190 MeV



Proton Energy Spectrum



Efficiency

	$\eta_E(\text{H}^+)^*$	$\#(\text{H}^+)^*$
>100MeV	1.1%	6.5×10^{10}
>50MeV	2.2%	1.9×10^{11}
>10MeV	4.2%	9.7×10^{11}

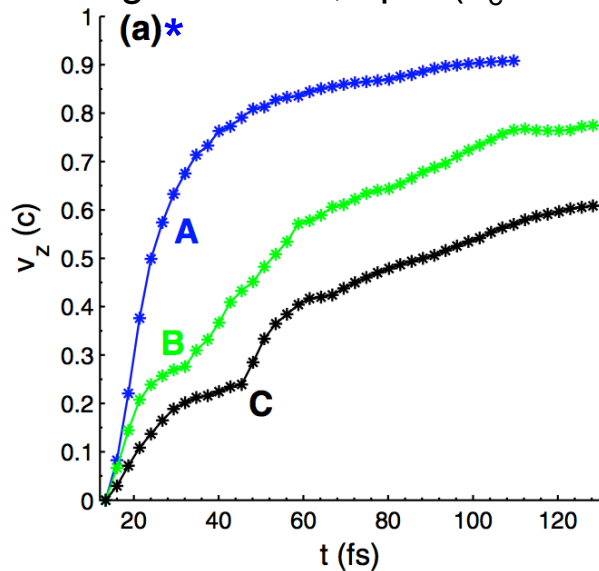
Monoenergetic protons with energy up to GeV can be obtained using higher intensity laser

*Qiao et al. (PRL 2009):

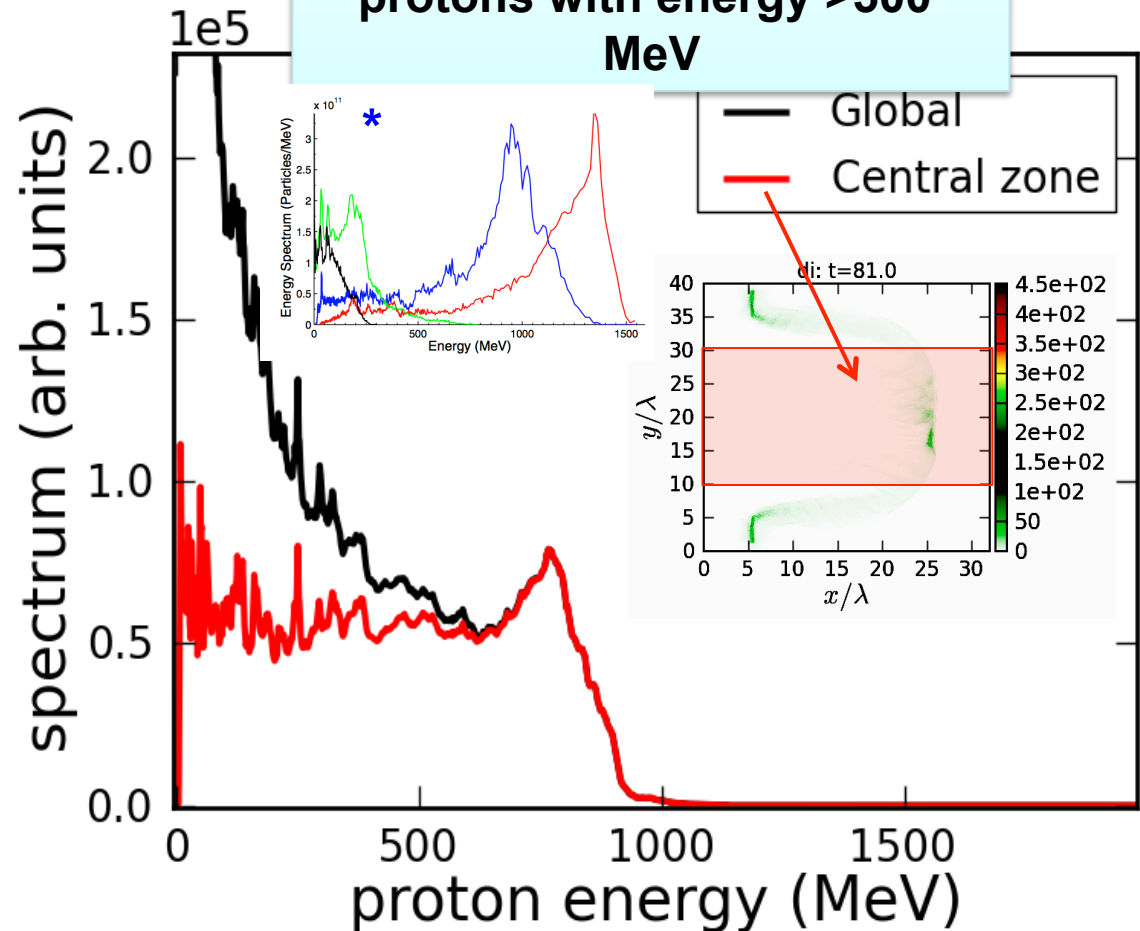
⇒ **Peak intensity: 1.89×10^{22} W/cm²**, Circ. Pol., 4th order super-Gaussian in space ($r=10\mu\text{m}$), Gaussian in time ($t=38$ fs)

⇒ $E_{\text{laser}} \sim 3.5$ kJ

■ Target: solid H, $1\mu\text{m}$ ($n_e=30 n_{\text{cr}}$)



19% of laser energy goes into 6×10^{12} forward directed protons with energy >500 MeV

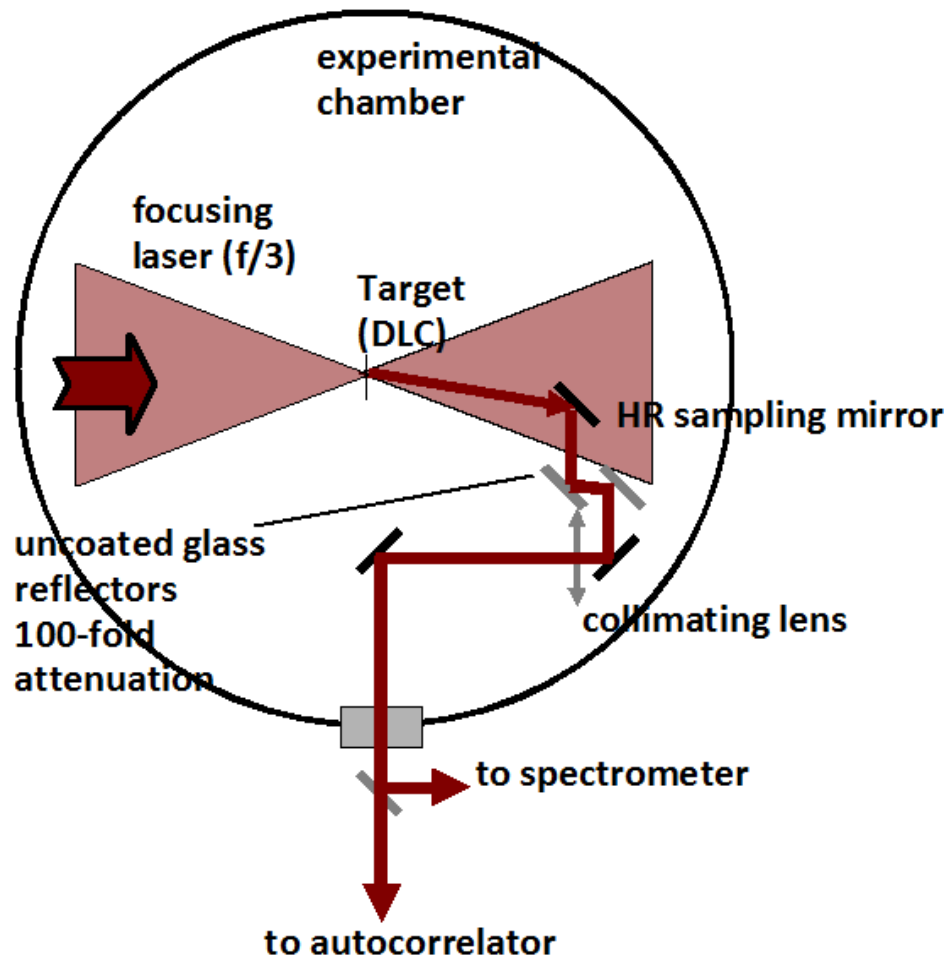


Nanotargets can...

... Perform laser pulse shaping

- Pulse shortening – 2nd order autocorrelation measurements
- Pulse shaping – simulations and diagnostic development

Experimental Setup

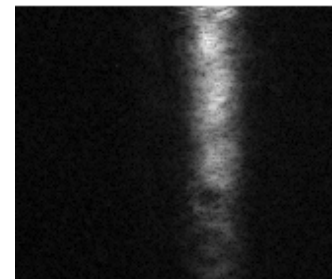


Glass reflectors attenuate intensity to avoid non-linear effects

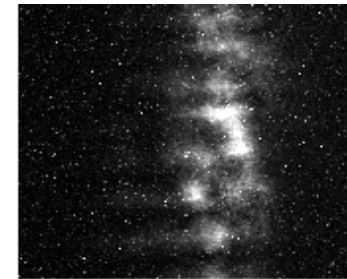
Spectra not analyzed yet

Small sampling ($\Phi 1\text{cm}$ of $\Phi 10\text{cm}$) ought guarantee uniformity (Fourier optics relates focal plane to spatial frequency in near field as $f_x = x/\lambda f_{\text{lens}}$. Frequencies $>1\text{ cm}^{-1}$ occur at focal positions $>100\text{ }\mu\text{m}$)

Non-uniformity induced by near-field diffraction from small aperture of sampling



front-end



target-side (no significant Target attenuation)

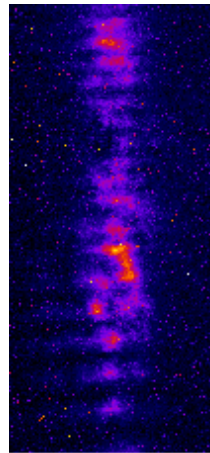
Results show shortening

Post-target

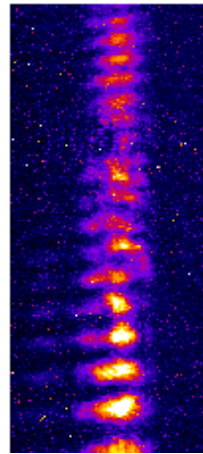
Autocorrelation images
(I/P > 520 fs for all)



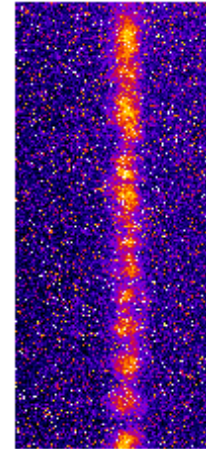
THINNER FOILS
no change



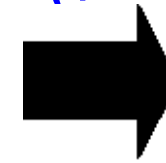
5 nm DLC



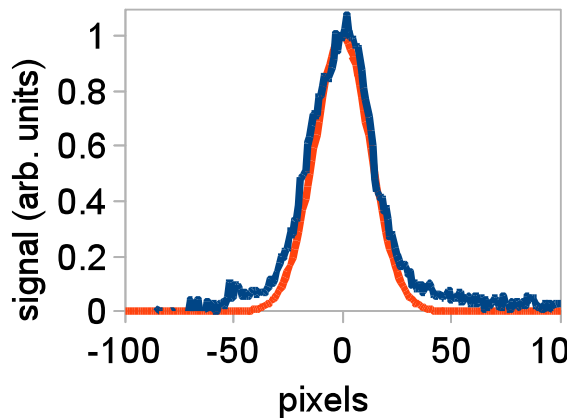
17 nm



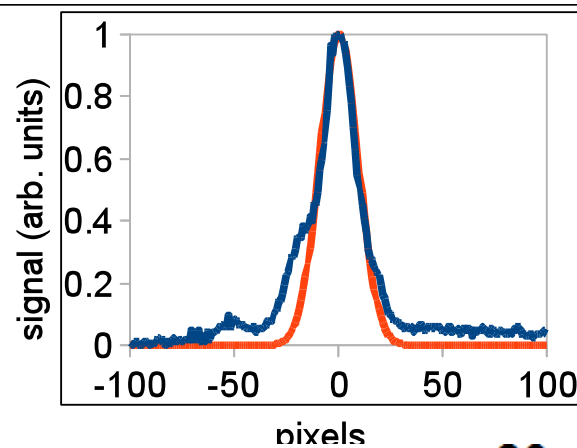
20 nm



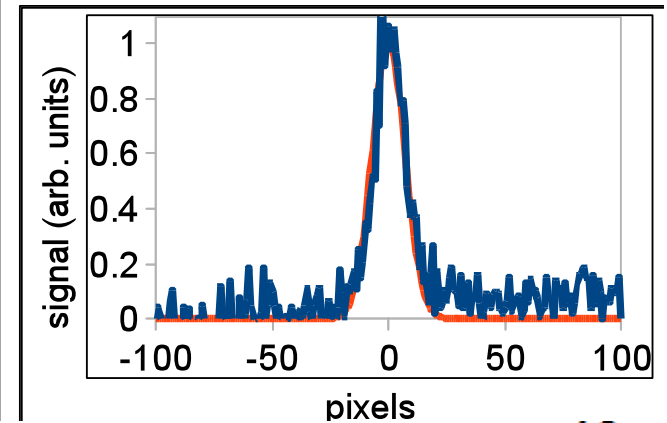
THICKER FOILS
Signal weakens and
below detection
at 300 nm



FWHM (pixels) 36
Duration (fs) 585



20
321

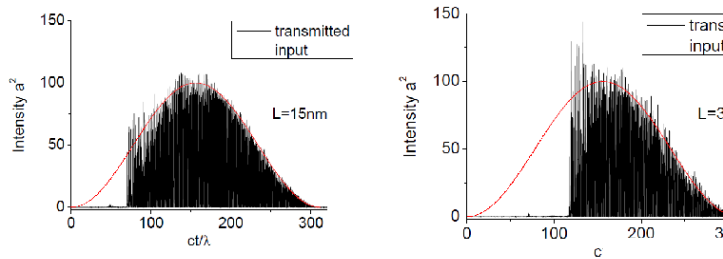


12
195

New single shot FROG works and will be able to measure truncated gaussian

Use nanofoil as temporal pulseshaper: simulations observe steep pulse cutoff

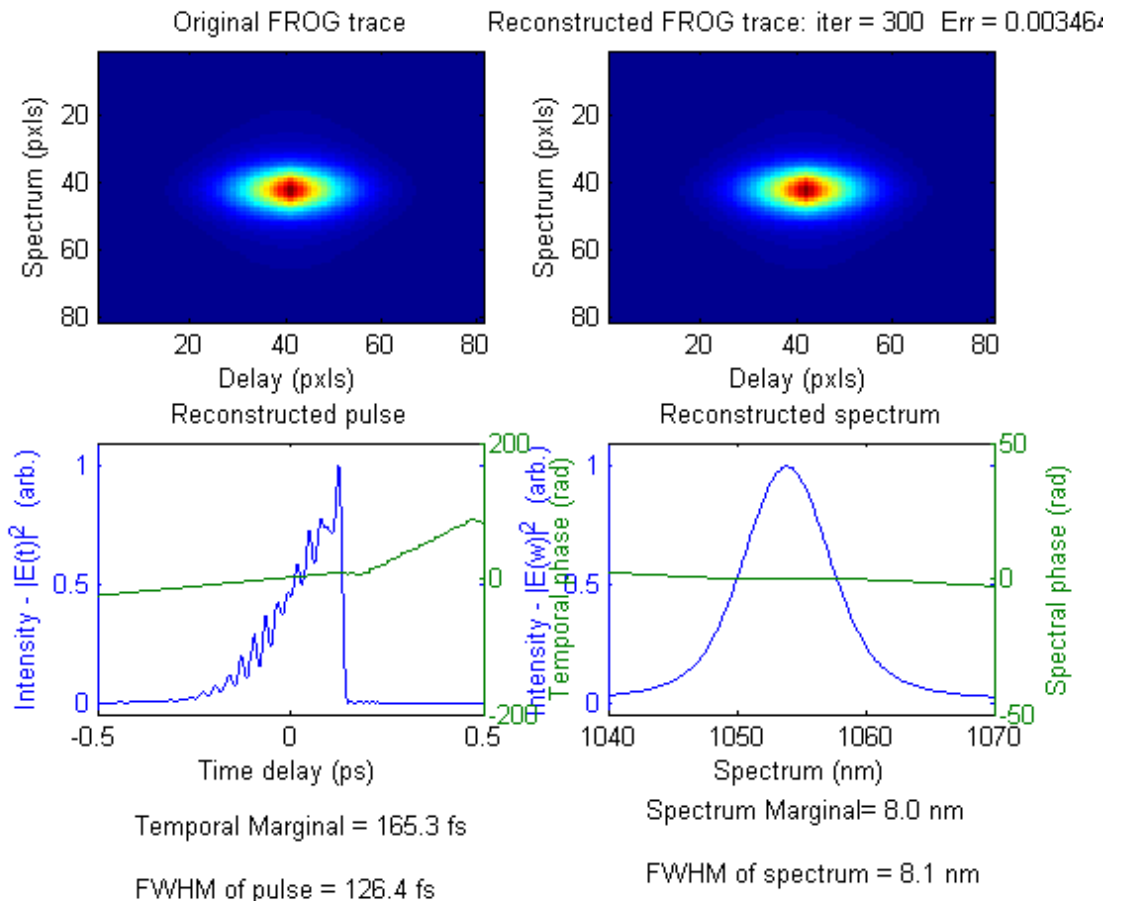
$a=10$, pulse FWHM=520fs, $n_e=700$ nc



1D VPIC simulation (Hui-Chun Wu)
 $a=10$, FWHM=520 fs,
 $n_e=700$ nc

Truncated Gaussian used to create FROG

Full Gaussian used to create FROG



Retrieved via algorithm

Nanotargets can be a coherent x-ray source

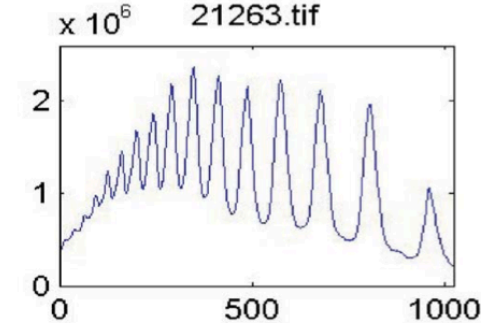
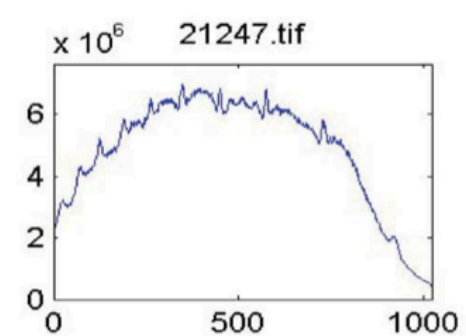
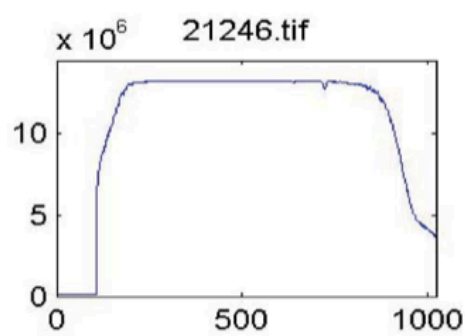
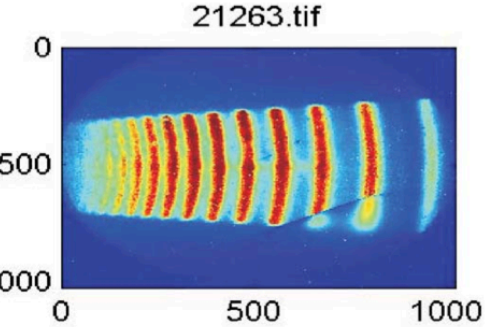
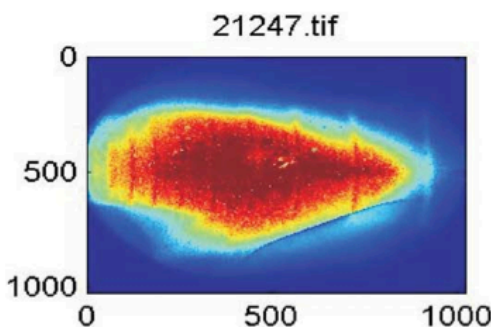
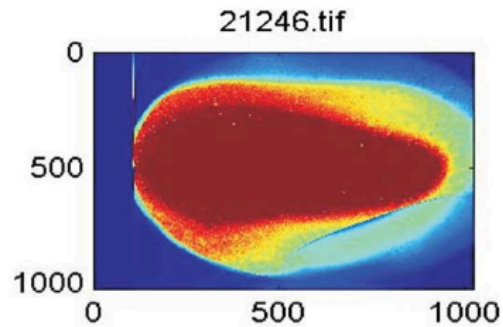
- Relativistic High Harmonics
- Relativistic Electron Mirrors
- Coherent Thomson Scattering

Forward relativistic high harmonic generation from nano-DLC targets

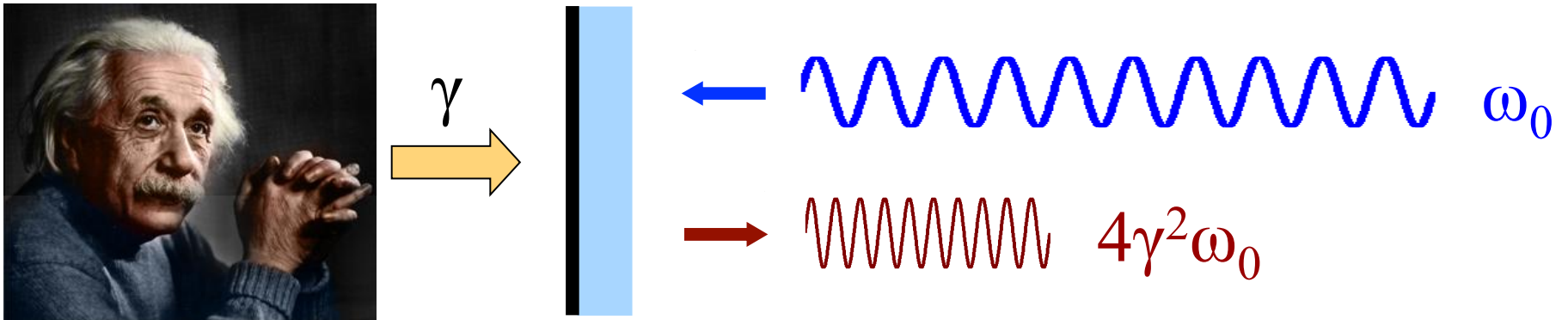
Date: 05.09.09 **Shot:** 21246
Laser: TRIDENT LANL
Energy: missing **Polarization:** LP
PM: no PM
Target: 10 nm DLC
Spectrometer: ACTON VM-502
Focus: 2"Quartz, 45 deg., rcc $\frac{1}{2}$ m
 λ_c : 100nm **MCP,Phos.:** -1/2 kV
Other:

Date: 05.09.09 **Shot:** 21247
Laser: TRIDENT LANL
Energy: 80J **Polarization:** LP
PM: no PM
Target: 50nm DLC
Spectrometer: ACTON VM-502
Focus: 2"Quartz, 45 deg., rcc $\frac{1}{2}$ m
 λ_c : 100nm **MCP,Phos.:** -1/2 kV
Other:

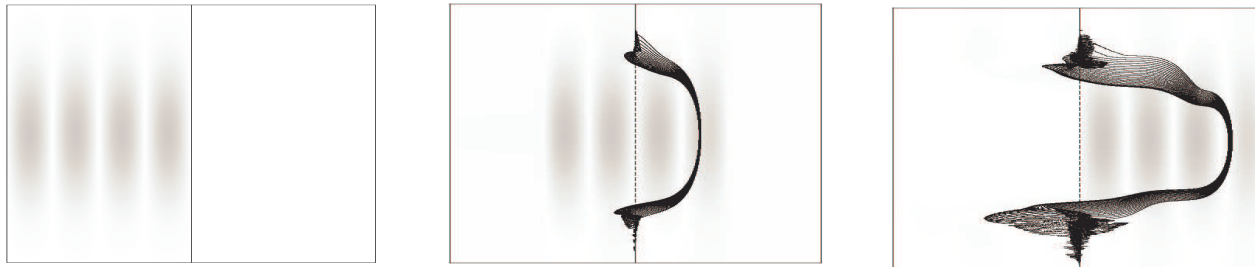
Date: 14.09.09 **Shot:** 21263
Laser: TRIDENT LANL
Energy: missing **Polarization:** LP
PM: no PM
Target: 200nm DLC
Spectrometer: ACTON VM-502
Focus: 2"Quartz, 45 deg., rcc $\frac{1}{2}$ m
 λ_c : 80nm **MCP,Phos.:** -1/2 kV
Other:



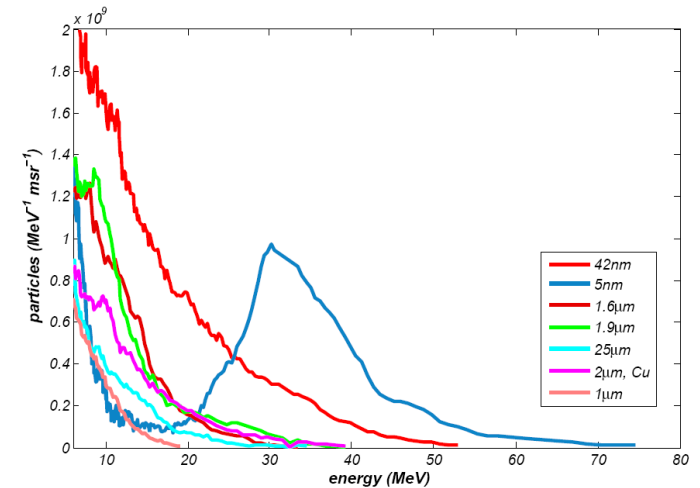
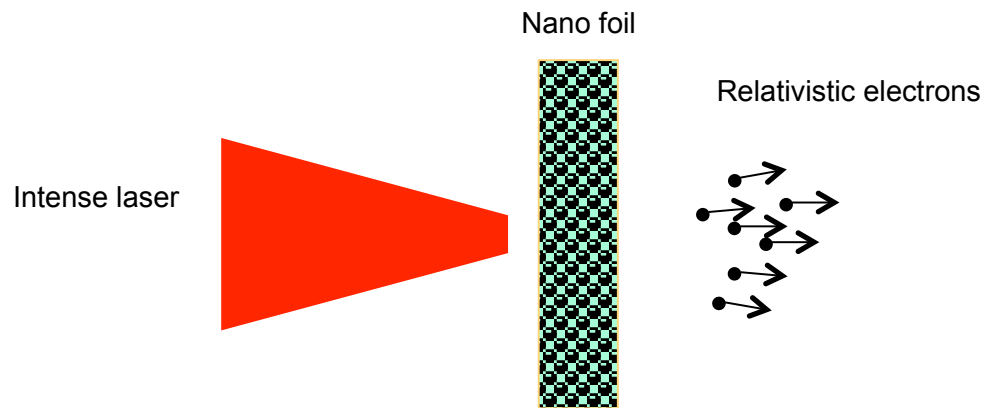
Coherent Thomson Scattering off a laser-driven Relativistic Electron Mirror



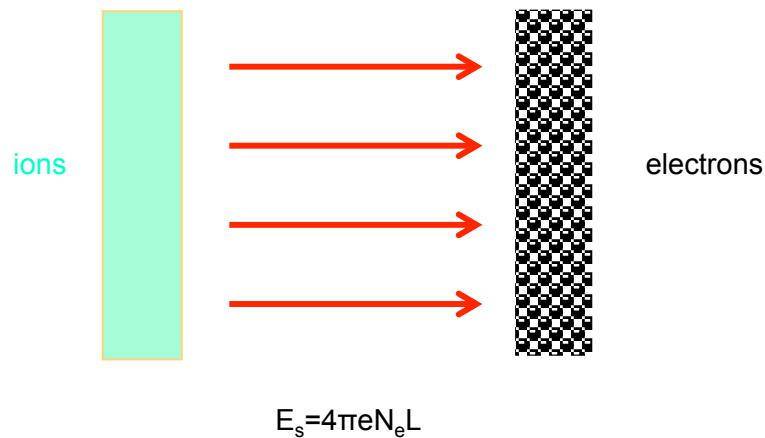
A. Einstein, 1905.



For very thin foils laser force can overcome restoring force of ions: Electron break-out observed on Trident!



electron mirror ?



Parameter region :

$$E_L \gg E_s$$

$$a_0 \gg \epsilon_0 = 2\pi(N_e/N_c)(L/\lambda_0)$$

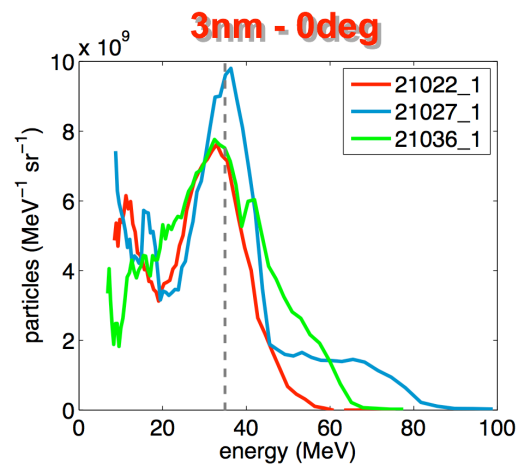
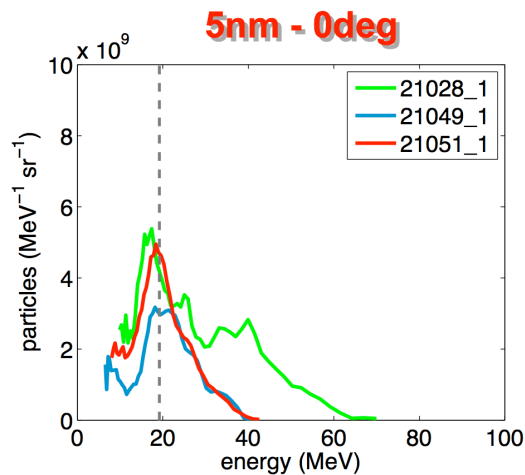
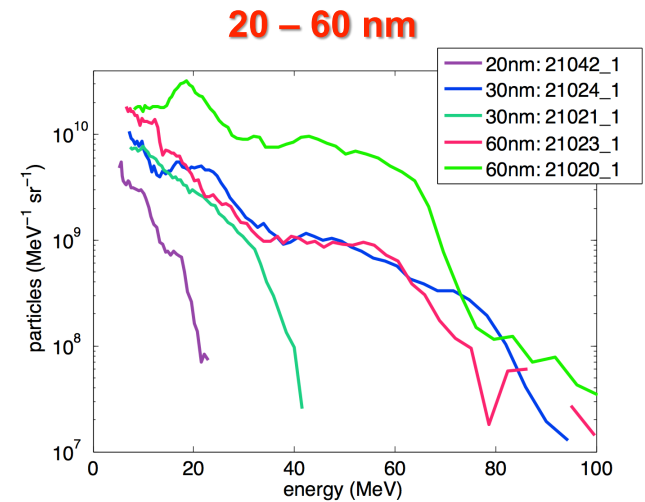
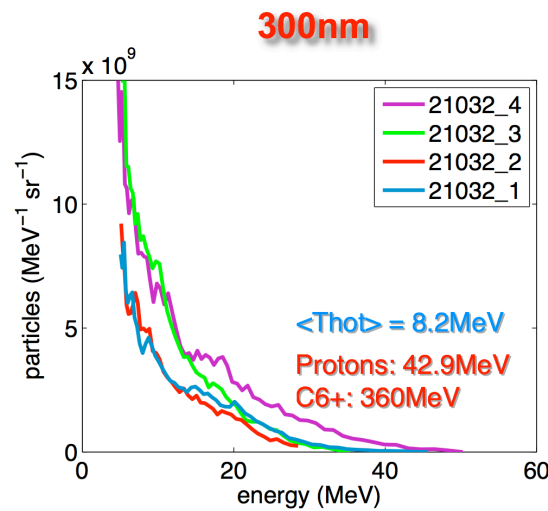
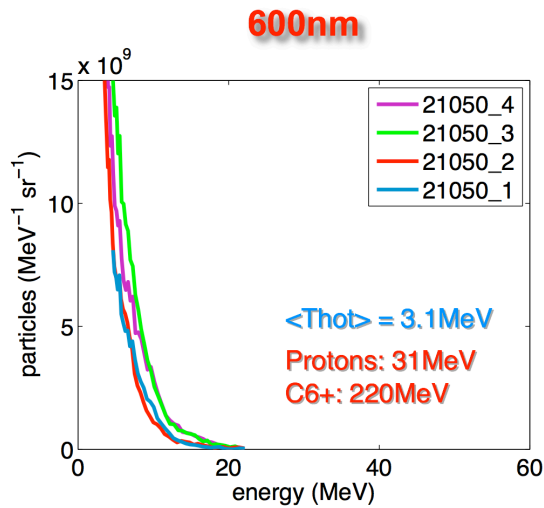
(Norm.)

$$I_0 = a_0^2 \times 1.37e18 \text{ W/cm}^2 /$$

$$(\lambda_{\mu\text{m}})^2$$

$$a_{0,\text{Trident}} = 20, \epsilon_{0,\text{Graphene}} = 1$$

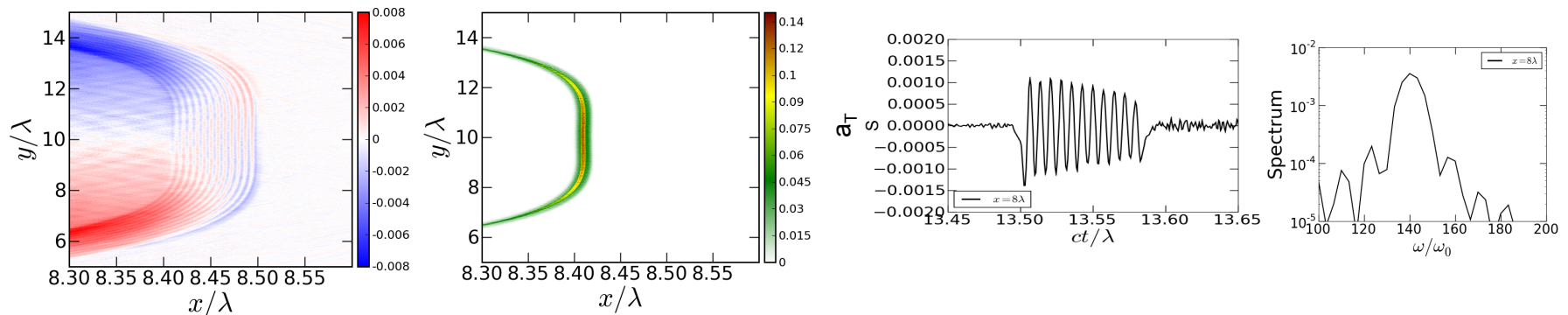
Electron temperature rises with decreasing foil thickness until spectra become mono-energetic



Typical laser parameters:

- $E = 80 - 100 \text{ J}$ on target
- $\tau = 500 - 600 \text{ fs}$ FWHM
- contrast $\sim 10^{-12}$ @ 0.5 ns
- contrast $\sim 10^{-7}$ @ 5 ps
- $50\% E_{\text{laser}}$ in $r_{\text{ave}} = 3.5 \mu\text{m}$
- $10\% E_{\text{laser}}$ in $r_{\text{ave}} = 1 \mu\text{m}$
- $I_{3.5\mu\text{m}} = 2 \times 10^{20} \text{ W/cm}^2$
- $a_{\text{ave}} = 13$
- $I_{3.5\mu\text{m}} = 5 \times 10^{20} \text{ W/cm}^2$
- $a_{\text{ave}} = 20$

Expected Performance:



- $a_0=3.5$: $\lambda_{ts}=5.7\text{nm}$ (217eV) and duration 267as.
- coherence w/o compression:
 - $N_e \sim 10^{24}$ e-/cc $\Rightarrow \lambda_{ts} = 10^{-8}$ cm $\Rightarrow E_{\text{photon}} = 12$ keV
 - $\Rightarrow \gamma = 55, a_0 = 10, I = 10^{20}$ W/cm²
- simulations observe 50x compression of e- layer:
 - $N_e \sim 5 \times 10^{25}$ e-/cc $\Rightarrow \lambda_{ts} = 2.7 \times 10^{-9}$ cm $\Rightarrow E_{\text{photon}} = 45$ keV

Summary

- Combination of nm-scale targets and ultrahigh contrast pulses leads to a paradigm shift in relativistic laser-matter interaction.
- We have demonstrated the production and integration of nm-scale targets and ultraclean pulses.
- This enables new particle acceleration mechanisms (BOA, PSA/ALPA):
 - IFI energies (500 MeV) reached with modest laser
 - IFI efficiencies (>10%) achieved: 80 MeV C6+ @ 11% CE with 0.7J Ti:Sa laser
 - IFI spectrum ($\Delta E < 20\%$) demonstrated: at Trident with 5nm targets and circular polarization.
 - Laser-pulse shortening and pulse shaping beyond Fourier-transform limit
 - Forward directed HHG from nano-DLC targets
 - Demonstration of quasi-monoenergetic electrons and electron break-out for sub-10nm targets important first step towards Relativistic Electron Mirrors (REM)