# Optimization and Limitation of Table-top Free-Electron Lasers

Sven Reiche Workshop "Frontier in Intense Laser-Matter Interaction" March 3<sup>rd</sup> 2010, Munich

#### Free-Electron Lasers



## **General Optimization**

#### **Electron Beam Parameters**

Almost everything scales with the FEL parameter:

$$\rho = \frac{1}{2\gamma} \left[ \left( \frac{\lambda_u f_c a_w}{2\pi} \right)^2 \frac{I}{I_A \sigma_x \sigma_y} \right]^2$$

- Assuming a round beam with similar emittance and symmetric focusing, the FEL parameter scales as  $(1/\epsilon_n)^{1/3}$ .
- The characteristic length is the power gain length:

$$L_g = \frac{\lambda_u}{4\sqrt{3}\pi\rho}$$
 (Ideal Case)

 Increasing the current and/or reducing the emittances increase the performance and reduce the overall required length of the FEL.

### Degrading Effects

	Energy Spread	Space Charge
Effect	Smears out micro bunching	Work against Coulomb field when bunching
Negligible when	$\frac{\sigma_{\gamma}}{\gamma} << \rho$	$k_p = \sqrt{\frac{4\pi I}{\gamma I_A A}} << 2k_u \rho \gamma$
Correction to FEL parameter	$1 - \left(\frac{\sigma_{\gamma}}{\rho\gamma}\right)^2$	$1 - \frac{1}{3} \left( \frac{k_p}{2k_u \rho \gamma} \right)^2$
Limitation	At lower energies	At lower energies or very high currents

#### Saturation Power and Brilliance

• Saturation power:

$$P_{sat} = 1.6\rho P_{beam} = 1.6\frac{mc^2}{e}\rho\gamma I$$

- For a given wavelength it favors a higher beam energy and thus a longer undulator period.
- Peak Brilliance:

$$Brill. = \frac{N_{photons}}{2\pi^{3}\sigma_{t}(\sigma_{\omega} / \omega)\sigma_{x}\sigma_{x'}\sigma_{y}\sigma_{y'}}$$

• At saturation:

$$N_{phot} = P_{sat} \sqrt{2\pi} \sigma_t / \hbar \omega$$
  

$$\sigma_{\omega} / \omega = 2\rho$$
  

$$\sigma_x \sigma_{x'} = c / 2\omega \quad \text{(transverse coherence)}$$

$$Brill. \approx \frac{0.8}{e \lambda_e \lambda} \gamma I$$



#### **Undulator Optimization**

• Increasing aw increases the coupling between electron field and radiation field  $(a_w/\gamma)$ , however at very large values the energy has to increase to preserve the resonant wavelength ( $\gamma \sim a_w$ ).

$$a_w \ge 0.8$$

- A small undulator period reduces the overall size of the FEL, however a strong undulator field would require a fraction of the period as the gap:
  - Strong impact of undulator wakefields (~gap<sup>.2</sup>)
  - Strong focusing to allow full transmission through the smaller aperture

For very short periods an RF/laser wiggler is more feasible.

#### **Optimizing the Focusing**

• Transverse betatron motion delays the particle with respect to the on-axis particle. The delay scales with the focusing strength:



- RMS averaging over the beam the action variables are replaced by the normalized emittance (and the energy deviation by the energy spread)
- Emittance effects are not disrupting the FEL performance if the condition is fulfill:

 $\frac{\lambda_u \varepsilon_n}{2 \lambda R_M} << \rho$ 

#### Optimizing the Focusing II

- Decreasing the  $\beta$ -function (increase focusing), increases the FEL parameter  $\rho$ .
- Too strong focusing enhances the emittance effect and increasing the FEL gain length.



#### Transverse Coherence (2D FEL Theory)

• Diffraction Parameter:



• Assuming electron size as radiation source size:



#### Transverse Coherence (2D FEL Theory)

• Growth rates for FEL eigenmodes (r,φ-decomposition):



#### SASE and Partial Coherence

- Spontaneous radiation as seed couples to many modes.
- Mode content visible in fluctuation of instantaneous power:

$$\zeta = \frac{\left\langle \left( P - \langle P \rangle \right)^2 \right\rangle}{\left\langle P \right\rangle^2} = \frac{1}{M_T}$$

 Agrees (surprisingly) well with standard definition of coherence:

$$\zeta = \frac{\iint \left| \mu_{12}(r_1, r_2) \right|^2 \left\langle I(r_1) I(r_2) \right\rangle dr_1 dr_2}{\left[ \int \left\langle I(r) \right\rangle dr \right]^2}$$

with the mutual intensity function:

$$\mu_{12}(r_1, r_2) = \frac{\left\langle E(r_1) E^*(r_2) \right\rangle}{\left[ \left\langle \left| E(r_1) \right|^2 \right\rangle \left\langle \left| E(r_2) \right|^2 \right\rangle \right]^{\frac{1}{2}}}$$

$$\mu_{12}$$
=1 : full coherence between  $r_1$  and  $r_2$ 



Case Study: High Current Beam Soft X-ray FEL Micro-Undulator

#### **Electron Beam Parameters**

Energy	1 GeV
Current	10 kA
Emittance	1 mm mrad
Energy spread	5 MeV

• Minimum wavelength (emittance constraint):

#### $\lambda > 1 \text{ nm}$

• Micro undulator:

 $a_w = 0.7, \lambda_u = 5 \text{ mm}$ 

• Field and Gap Estimate (planar hybrid undulator):

B=2.1 T, g = 0.5 mm

(gap can be increased with kryogenic undulators)

#### FEL Performance (Ming Xie Model)

• Strongly effected by the energy spread and sub-sequentially by the emittance. The simple 1D FEL parameter is about 2% but reduces by factor of 20 in the 3D model.

Effective FEL Parameter	0.1%
Gain Length	0.23 m
Saturation Length	4.8 m
Saturation Power	9 GW
Bandwidth	0.23 % FWHM

- Requires 5 m long micro undulator.
- Performance goes down (less power, reduced coherence) for:
  - Shorter wavelength / higher energy
  - Larger energy spread

Case Study: Ultralow Emittance Beam Hard X-ray FEL Laser Wiggler

#### The Advantage of Laser Wigglers

- Counter-propagating laser fields have the same impact on the electron beam as an undulator field. (in electron rest frame, the undulator field becomes an EM wave).
- The period can be reduced significantly while keeping a strong field (a ~1)
- Avoids the boundary problems of a magnetic undulator, which are:
  - Physical aperture between the poles
  - Strong wakefields within the undulator.
- Tunability of the *a*-value with the pulse energy.

#### Requirement for the Laser Field.

• Minimum interaction volume:



- Within the interaction volume the field stability over the interaction time L<sub>b</sub>/c has to be:
- The absolute minimum Rayleigh length is (Mode stacking):
  - Fundamental transverse mode only:

$$\frac{\Delta a}{a} \le \rho$$



#### Breaking the Angstrom Barrier...

- Goal: 1 Ångstrom radiation, using Ti:Saph laser wiggler
- Electron beam:  $E = a \cdot 50 \text{ MeV}, \epsilon_n = a \cdot 10^{.9} \text{ m}$  (!!!)
- Beam source is most likely limited to current < 100 A.
- Beta-function about 10 cm for tight spot of  $r_0=1 \ \mu m$ .

Huge penalties

$$\rho = \frac{1}{2\gamma} \left[ \left( \frac{\lambda_u f_c a_w}{2\pi} \right)^2 \frac{I}{I_A \sigma_x \sigma_y} \right]^{\frac{1}{3}} \approx 10^{-4}$$

Additional problems with coherence and energy spread.

#### Making Laser (based) Wiggler Realizable...

- Relaxed condition for longer period (~100  $\mu$ m):
  - THz Radiation (doubtful)
  - Laser Beatwave Plasma Wiggler (channeled, but very sensitive to spot size, requires even lower emittance)
- Work with higher field strength (a>>1) to relax demands on electron beam emittance and allow higher beam currents.
- Active control of FEL parameter ρ:
  - Too small → field stability/energy spread requirements
  - Too high → reduced coherence

#### Summery

- Beam emittance defines the achievable wavelength:
  - → Small emittance value
  - → Lower beam energy
  - Shorter period length
  - → Shorter saturation length.
- Coherence seems to be the most limiting factor when going to shorter wavelength.
- With current e-sources, 1 nm seems reasonable but 1 Å requires significant improvement in beam quality
- Saturation power and peak brightness will be lower than Xray FEL facilities (XFEL, LCLS)
- Micro undulators seem to be the preferred choice, laser wiggler have many technical issues, which needs to be resolved.